

# Monolithic Newton multigrid FEM techniques for nonlinear problems with special emphasis on viscoelastic fluids

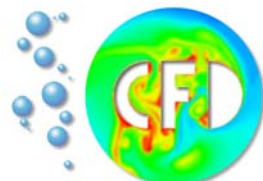
H. Damanik, J. Hron, A. Ouazzi, S. Turek

Institut für Angewandte Mathematik, LS III, TU Dortmund

<http://www.mathematik.tu-dortmund.de/lisiii>

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University of Leicester,  
Leicester, UK



- Generalized Navier-Stokes equations

$$\rho \left( \frac{\partial}{\partial t} + u \cdot \nabla \right) u = -\nabla p + \nabla \cdot \sigma, \quad \nabla \cdot u = 0,$$

$$\rho c_p \left( \frac{\partial}{\partial t} + u \cdot \nabla \right) \Theta = k_1 \nabla^2 \Theta + k_2 D : D,$$

$$D(u) = \frac{1}{2} \left( \nabla u + (\nabla u)^T \right),$$

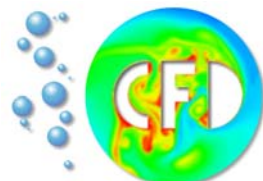
$$\sigma = \sigma_s + \sigma_p.$$

- Viscous stress

$$\sigma_s = 2 \eta_s (\dot{\gamma}, \Theta, p) D, \quad \dot{\gamma} = \sqrt{\text{tr}(D(u)^2)}.$$

- Elastic stress

$$\sigma_p + We \frac{\delta_a \sigma_p}{\delta t} + G(\sigma_p, D) + H(\sigma_p) = 2 \eta_p D(u).$$



- **Viscous stress**

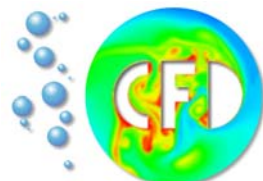
$$\sigma_s = 2 \eta_s(\dot{\gamma}, \Theta, p) D, \dot{\gamma} = \sqrt{\text{tr}(D(u)^2)}.$$

- **Power law model**

$$\eta_s(\dot{\gamma}, \Theta, p) = \eta_0 (\varepsilon + \dot{\gamma}^2)^{\left(\frac{r}{2}-1\right)}, (\eta_0 > 0, r > 1).$$

- **Generalized Cross model**

$$\eta_s(\dot{\gamma}, \Theta, p) = \eta_\infty + \frac{(\eta_0 - \eta_\infty)}{(1 + (\lambda \dot{\gamma})^{r_1})^r} \exp\left(\alpha p + \left(a_1 + \frac{a_2}{a_3 + \Theta}\right)\right),$$
$$(\eta_0 > \eta_\infty \geq 0, r > 1, \lambda > 0).$$



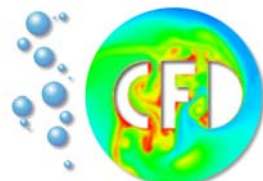
- Elastic stress

$$\sigma_p + We \frac{\delta_a \sigma_p}{\delta t} + G(\sigma_p, D) + H(\sigma_p) = 2\eta_p D(u).$$

- Upper/Lower convective derivative

$$\frac{\delta_a \sigma}{\delta t} = \left( \frac{\partial}{\partial t} + u \cdot \nabla \right) \sigma + g_a(\sigma, \nabla u)$$

$$g_a(\sigma, \nabla u) = \frac{1-a}{2} (\sigma \nabla u + (\nabla u)^T \sigma) - \frac{1+a}{2} (\nabla u \sigma + \sigma (\nabla u)^T), \quad a = \pm 1$$



- **Generalized differential constitutive model**

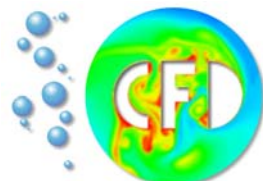
$$\sigma_p + We \frac{\delta_a \sigma_p}{\delta t} + G(\sigma_p, D) + H(\sigma_p) = 2\eta_p D(u).$$

- **Oldroyd-B**  $G = 0, H = 0, a = 1.$

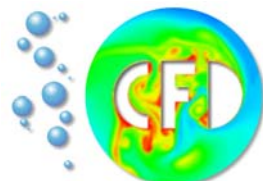
- **Giesekus**  $G = 0, H = \alpha \sigma_p^2, a = 1.$

- **Phan-Thien and Tanner**  
 $G = 0, H = [\exp(\alpha \text{tr}(\sigma_p)) - 1] \sigma_p.$

- **White and Metzner**  
 $G = \alpha (2D : D)^{\frac{1}{2}}, H = 0.$



- FEM techniques have to handle the following challenges points
  - Stable FE spaces for velocity/pressure and velocity/stress interpolation  $Q_2 / Q_2 / P_1^{disc}$  or  $\tilde{Q}_1 / \tilde{Q}_1 / P_0$  or the new  $\tilde{Q}_2 / \tilde{Q}_2 / P_1^{disc}$
  - Special treatment of the convective terms: **edge-oriented/interior penalty (EO-FEM), TVD/FCT**
  - High Weissenberg number problem (HWNP): **LCR**
- Solvers have to deal with different sources of nonlinearity
  - Nonlinearity: **Newton method**
  - Strong coupling of equations: **monolithic multigrid approach**
- Complex geometries (and meshes)



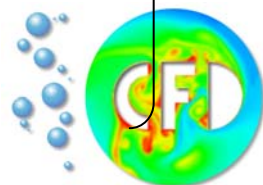
Old  $\rightarrow (u, p, \sigma_p)$

$$\rho \left( \frac{\partial}{\partial t} + u \cdot \nabla \right) u = -\nabla p + 2\nabla \cdot \eta_s D + \nabla \cdot \sigma_p, \nabla \cdot u = 0,$$
$$\sigma_p + We \frac{\delta_a \sigma_p}{\delta t} = 2\eta_p D,$$

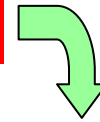
**Conformation tensor**  $\rightarrow (u, p, \sigma_c)$  This tensor is positive definite by design !!

Replace  $\sigma_p$  in (1) with  $\sigma_p = \frac{\eta_p}{We} (\sigma_c - I)$   $\rightarrow$  special discretization : TVD

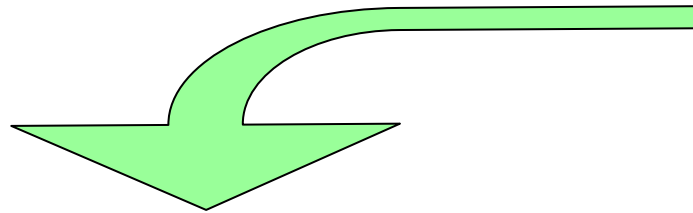
$$\rho \left( \frac{\partial}{\partial t} + u \cdot \nabla \right) u = -\nabla p + 2\nabla \cdot \eta_s D + \frac{1}{We} \nabla \cdot \eta_p \sigma_c, \nabla \cdot u = 0,$$
$$\frac{\delta_a \sigma_c}{\delta t} + \frac{1}{We} (\sigma_c - I) = 0,$$



$$\sigma_c(t) = \int_{-\infty}^t \frac{1}{We^2} \exp\left(\frac{-(t-s)}{We}\right) F(s,t) F(s,t)^T ds$$

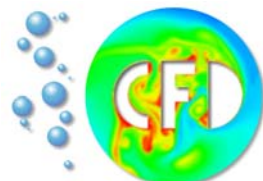


Positive by design,  
so we can take its logarithm



## 2 Observations:

- positive definite → special discretizations like FCT/TVD
- exponential behaviour → approximation by polynomials???



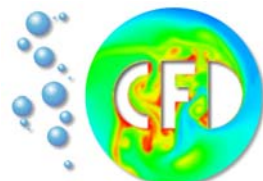


Replace  $\sigma_c$  in (2) with  $\sigma_c = e^{\sigma_{lc}}$

$$\left. \begin{aligned} \rho \left( \frac{\partial}{\partial t} + u \cdot \nabla \right) u &= -\nabla p + 2\nabla \cdot \eta_s D + \frac{1}{We} \nabla \cdot \eta_p e^{\sigma_{lc}}, \quad \nabla \cdot u = 0, \\ \frac{\delta_a e^{\sigma_{lc}}}{\delta t} + \frac{1}{We} (e^{\sigma_{lc}} - I) &= 0, \end{aligned} \right\} (3)$$

Gradient of exponential of  $e^{\sigma_{lc}} \rightarrow ???$

Solvers  $\rightarrow ???$



- **Experience:**

- **Stresses grow exponentially**
- **Conformation tensor is positive by design**

- **Idea**

- **Take the logarithm** as a new variable  $\sigma_{lc} = \log \sigma_c$  using the eigenvalue decomposition

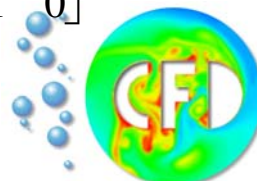
$$\sigma_{lc} = R \log(\lambda_{\sigma_c}) R^T$$

- **Decompose the velocity gradient inside the stretching part**

$$\nabla u = \Xi + Y + \bar{\Xi} \tau^{-1}$$

$$\nabla_{\sigma_c} u = R^T (\nabla u) R, Y(\nabla_{\sigma_c} u) = R [\text{diag}(\nabla_{\sigma_c} u)] R^T, \Xi = R[\omega \chi] R^T, \bar{\Xi} = R[\bar{\omega} \chi] R^T, \chi = \begin{bmatrix} 0 & 1 \\ -1 & 0 \end{bmatrix}$$

$$\omega = \lambda_2 [\nabla_{\sigma_c} u]_{12} + \lambda_1 [\nabla_{\sigma_c} u]_{21} / (\lambda_2 - \lambda_1), \bar{\omega} = [\nabla_{\sigma_c} u]_{21} + [\nabla_{\sigma_c} u]_{12} / (\lambda_2^{-1} - \lambda_1^{-1}).$$



- **Generalized Newtonian (VP)**

$$\rho \left( \frac{\partial}{\partial t} + u \cdot \nabla \right) u = -\nabla p + \nabla \cdot (2\eta_s(\dot{\gamma}, \Theta, p) D(u)) + \frac{1}{We} \nabla \cdot \eta_p e^{\sigma_{lc}}, \quad \nabla \cdot u = 0,$$

- **Non-isothermal effect (T)**

$$\rho c_p \left( \frac{\partial}{\partial t} + u \cdot \nabla \right) \Theta = k_1 \nabla^2 \Theta + k_2 D : D,$$

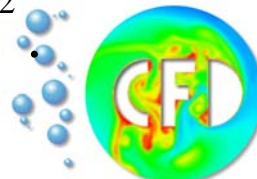
- **Stress equation (S)**

$$\left( \frac{\partial}{\partial t} + u \cdot \nabla \right) \sigma_{lc} - (\Xi \sigma_{lc} - \sigma_{lc} \Xi) - 2Y = \frac{1}{We} f(\sigma_{lc}).$$

- **Oldroyd-B model**

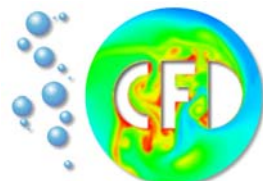
$$f(\sigma_{lc}) = (e^{-\sigma_{lc}} - I)$$

- **Giesekus model**  $f(\sigma_{lc}) = (e^{-\sigma_{lc}} - I) - \alpha e^{\sigma_{lc}} (e^{-\sigma_{lc}} - I)^2$



- High order  $Q_2 / Q_2 / P_1^{disc}$  for velocity-stress-pressure
  - Advantages:
    - Inf-sup stable for velocity and pressure
    - High order: good for accuracy
    - Discontinuous pressure: good for solver
  - Disadvantages:
    - Stabilization for same spaces for stress-velocity
    - a single d.o.f. belongs to four elements (in 2D)

**Compatibility condition between the stress and velocity spaces via EO-FEM !**

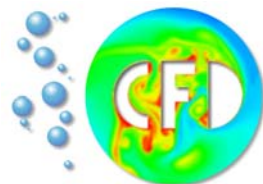


- **Edge-oriented stabilization for**
  - Same finite element interpolation velocity and stress
  - convective dominated problem

$$\langle J_{\tilde{u}} \tilde{u}, \tilde{v} \rangle = \sum_{\text{edge} E} \max(\gamma_u \eta_p h_E, \gamma_{\tilde{u}} h_E^2) \int_E [\nabla \tilde{u}] [\nabla \tilde{v}] ds$$

with  $\tilde{u} = (u, \sigma, \Theta)$ ,  $\tilde{v} = (v, \tau, \Phi)$  and  $[\nabla \tilde{u}] [\nabla \tilde{v}] = \sum_i [\nabla \tilde{u}_i] [\nabla \tilde{v}_i]$

**Efficient Newton-type and multigrid solvers  
can be easily applied !**



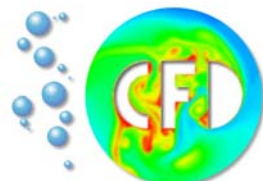
- **Damped Newton results in the solution of the form**

$$R(x) = 0, \quad x = (u, \sigma, \Theta, p)$$

$$x^{l+1} = x^l + \omega^l \left[ \frac{\partial R(x^l)}{\partial x} \right]^{-1} R(x^l)$$

- **Inexact Newton:** Jacobian is approximated using finite differences

$$\left[ \frac{\partial R(x^l)}{\partial x} \right]_{ij} \approx \frac{R_j(x + \varepsilon e_i) - R_j(x - \varepsilon e_i)}{2\varepsilon}$$



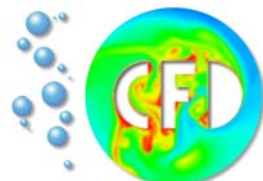
- The Jacobian matrix takes the form

$$\left[ \frac{\partial \mathcal{R}(x^n)}{\partial x} \right] = \begin{bmatrix} A & \tilde{B}^T \\ B & 0 \end{bmatrix}$$

- Generalized non-isothermal non-Newtonian problem

$$A = \begin{bmatrix} A_u & \tilde{C}^T & \tilde{H}^T \\ C & A_\sigma & 0 \\ H & 0 & A_\Theta \end{bmatrix}$$

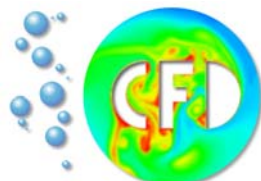
**Typical saddle point problem !**



- **Monolithic multigrid solver**
  - **Standard geometric multigrid approach**
  - **Full  $Q_2, P_1^{disc}$  restrictions and prolongations**
  - **Local MPSC via Vanka-like smoother**

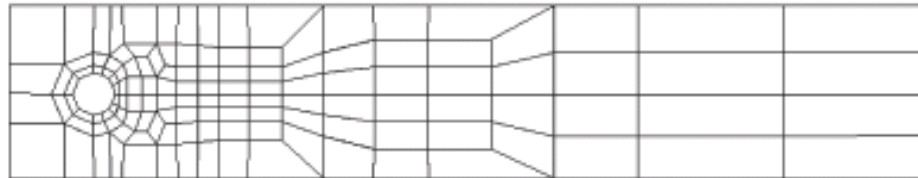
$$\begin{bmatrix} \tilde{u}^{l+1} \\ p^{l+1} \end{bmatrix} = \begin{bmatrix} \tilde{u}^l \\ p^l \end{bmatrix} + \omega^l \sum_{T \in \mathcal{T}_h} [A|_T]^{-1} \begin{bmatrix} R_u(\tilde{u}^l, p^l) \\ R_p(\tilde{u}^l, p^l) \end{bmatrix}_T$$

**Fully implicit Monolithic FEM-Multigrid Solver !**





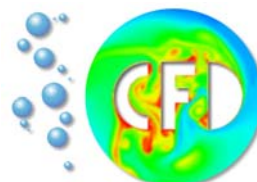
- Croarse grid and mesh information**



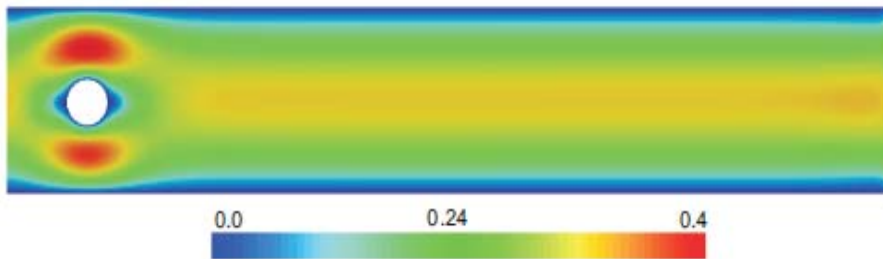
Level	NEL	NMT	NMP
1	520	572	1092
2	2080	2184	4264
3	8320	8528	16848
4	33280	33696	66976
5	133120	133952	267072

- n.o.f. for different problems**

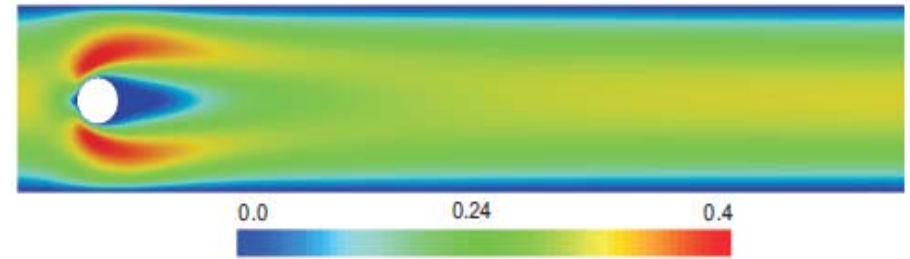
Level	n.o.f. u	n.o.f. p	n.o.f. T	n.o.f. S	VP	VPT	VSP
0	1144	390	572	1716	1534	2106	3250
1	4368	1560	2184	6552	5928	8112	12489
2	17056	6240	8528	25584	23296	31824	48880
3	67392	24960	33696	101088	92352	126048	193449
4	267904	99840	133952	401856	367744	501696	769600
5	1068288	399360	534144	1602432	1467648	2001792	3070080



## Stokes

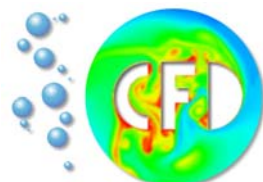


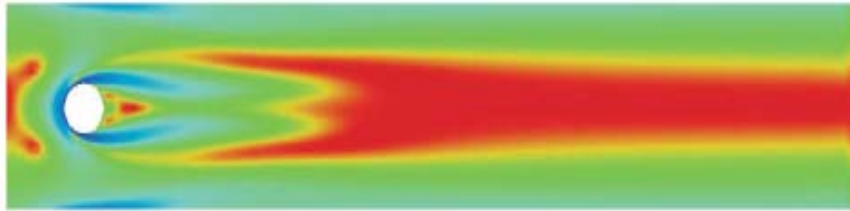
## Navier-Stokes Re=20



Level	Drag	Lift	N/L	Drag	Lift	N/L
1	3.112646	2.965870e-2	1/6	5.540999	9.447473e-3	5/2
2	3.134342	3.005275e-2	1/7	5.566928	1.046885e-2	5/2
3	3.140327	3.015909e-2	1/7	5.576088	1.056787e-2	5/2
4	3.141893	3.018665e-2	1/7	5.578652	1.060398e-2	5/2
5	3.142292	3.019366e-2	1/7	5.579313	1.061503e-2	5/2

**Level independent solver !**





$$\eta_s(\dot{\gamma}, \Theta, p) = \eta_0(\varepsilon + \dot{\gamma}^2)^{\frac{r-1}{2}}, (\eta_0 > 0, r > 1).$$

$\varepsilon = 10^{-2}, r = 1.5$

Level	Drag	Lift	N/L
2	3.26420	-0.01339	4/2
3	3.27728	-0.01341	3/2
4	3.27956	-0.01338	2/2
5	3.28007	-0.01337	2/2

$\varepsilon = 10^{-2}, r = 3$

Level	Drag	Lift	N/L
2	13.74280	0.35070	3/2
3	13.77355	0.34963	3/2
4	13.78220	0.35062	3/1
5	13.78445	0.35112	2/2

$\varepsilon = 10^{-4}, r = 1.5$

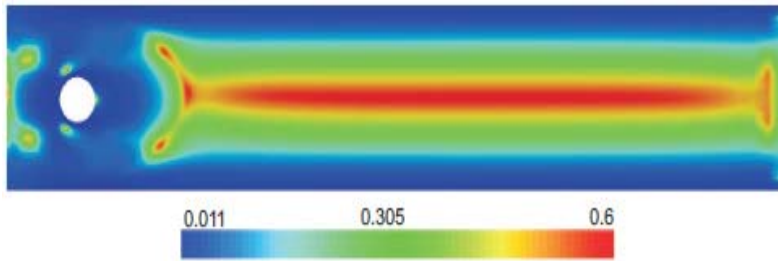
Level	Drag	Lift	N/L
2	3.26433	-0.01342	4/2
3	3.27739	-0.01342	3/2
4	3.27968	-0.01339	2/2
5	3.28019	-0.01338	2/2

$\varepsilon = 10^{-4}, r = 3$

Level	Drag	Lift	N/L
2	13.73800	0.35052	3/2
3	13.76875	0.34941	3/2
4	13.77740	0.35040	3/1
5	13.77970	0.35091	2/2

**Level and parameter independent solver !**





$$r = 1, r_1 = 1, \alpha = 0, \eta_0 = 10^{-2}$$

Level	Drag	Lift	N/L
2	6.31313	0.02478	3/1
3	6.32337	0.02504	3/2
4	6.32619	0.02509	3/2
5	6.32691	0.02510	2/2

$$r = 0, r_1 = 1, \alpha = 0.1, \eta_0 = 10^{-2}$$

Level	Drag	Lift	N/L
2	33.22763	0.81901	2/2
3	33.29026	0.82140	2/2
4	33.30657	0.82201	2/2
5	33.31069	0.82217	2/2

$$\eta_s(\dot{\gamma}, \Theta, p) = \eta_\infty + \frac{(\eta_0 - \eta_\infty)}{(1 + (\lambda \dot{\gamma})^{r_1})^r} \exp(\alpha p + (a_1 + \frac{a_2}{a_3 + \Theta})),$$

$$\eta_\infty = 10^{-3}, a_1 = a_2 = 0,$$

$$r = 0, r_1 = 1, \alpha = 0.1, \eta_0 = 10^{-1}$$

Level	Drag	Lift	N/L
2	534.29750	6.53247	3/2
3	535.48500	6.55813	3/3
4	535.77950	6.56464	3/3
5	535.84800	6.56621	2/2

$$r = 1, r_1 = 1, \alpha = 0.1, \eta_0 = 10^{-1}$$

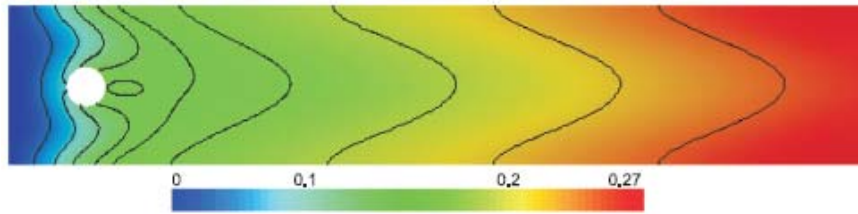
Level	Drag	Lift	N/L
2	15.16395	0.13886	4/3
3	15.18516	0.13963	4/3
4	15.19108	0.13978	4/3
5	15.19262	0.13982	3/3

**Level and model independent solver !**



# Cross model problem (VTP)

## Non heated cylinder



$$\eta_s(\dot{\gamma}, \Theta, p) = \eta_\infty + \frac{(\eta_0 - \eta_\infty)}{(1 + (\lambda \dot{\gamma})^{r_1})^r} \exp(\alpha p + (a_1 + \frac{a_2}{a_3 + \Theta})),$$

$$\eta_\infty = 10^{-3}, a_1 = 0, a_3 = 1, k_1 = k_2 = 10^{-2},$$

$$r = 0, r_1 = 1, \alpha = 0, \eta_0 = 10^{-2}, a_2 = 0.$$

Level	Drag	Lift	N/L
2	74.29465	1.31636	2/2
3	74.43290	1.32009	2/2
4	74.46910	1.32105	2/2
5	74.47830	1.32129	2/2

$$r = 0.1, r_1 = 1, \alpha = 0, \eta_0 = 10^{-2}, a_2 = 1.$$

Level	Drag	Lift	N/L
2	53.78930	1.05488	2/2
3	53.88590	1.05770	3/2
4	53.91125	1.05844	2/2
5	53.91770	1.05863	2/2

$$r = 0.1, r_1 = 1, \alpha = 0, \eta_0 = 10^{-1}, a_2 = 0.$$

Level	Drag	Lift	N/L
2	5579.355	54.79350	3/2
3	5589.415	54.96815	3/3
4	5592.050	55.01415	3/3
5	5592.725	55.02585	3/2

$$r = 0.1, r_1 = 1, \alpha = 10^{-3}, \eta_0 = 10^{-1}, a_2 = 1.$$

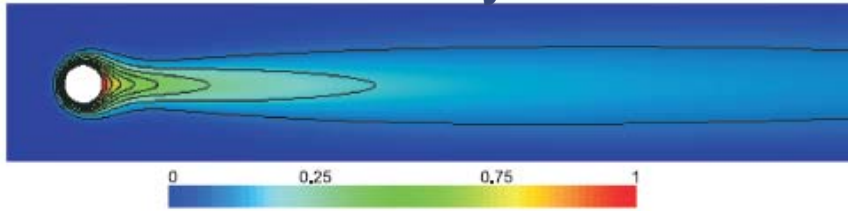
Level	Drag	Lift	N/L
2	6005.265	59.75125	3/2
3	6016.220	59.94455	3/2
4	6019.075	59.99535	3/2
5	6019.795	60.00820	3/2

**Level and model independent solver !**



# Cross model problem (VTP)

## Heated cylinder



$$\eta_0 = 10^{-2}, k_1 = k_2 = 10^{-2}.$$

Level	Drag	Lift	N/L
2	45.26969	0.90303	3/2
3	45.35251	0.90563	3/2
4	45.37431	0.90632	2/2
5	45.37988	0.90649	2/2

$$\eta_0 = 10^{-1}, k_1 = k_2 = 10^{-2}.$$

Level	Drag	Lift	N/L
2	464.07865	4.90752	2/2
3	464.93045	4.92313	3/2
4	465.15470	4.92724	2/2
5	465.21195	4.92828	2/2

$$\eta_s(\dot{\gamma}, \Theta, p) = \eta_\infty + \frac{(\eta_0 - \eta_\infty)}{(1 + (\lambda \dot{\gamma})^{r_1})^r} \exp(\alpha p + (a_1 + \frac{a_2}{a_3 + \Theta})),$$

$$\eta_\infty = 10^{-3}, r = 0.1, r_1 = 1, \alpha = 10^{-3}, a_1 = 0, a_2 = 1, a_3 = 1,$$

$$\eta_0 = 10^{-1}, k_1 = k_2 = 10^{-3}.$$

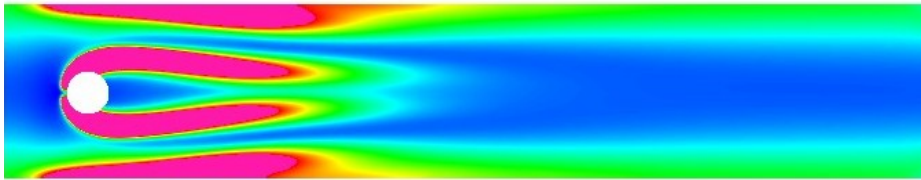
Level	Drag	Lift	N/L
2	512.7765	5.37640	3/2
3	513.7120	5.39301	3/3
4	513.9585	5.39743	3/3
5	514.0215	5.39856	3/2

$$\eta_0 = 1.0, k_1 = k_2 = 10^{-3}.$$

Level	Drag	Lift	N/L
2	5528.860	53.11685	3/2
3	5539.025	53.28790	3/2
4	5541.690	53.33335	3/2
5	5542.365	53.34490	3/2

**Level and model independent solver !**





$$\sigma_p + We \frac{\delta_a \sigma_p}{\delta t} + G(\sigma_p, D) + H(\sigma_p) = 2\eta_p D(u)$$

$We = 0.002$        $a = 1, G = 0, H = 0$

Level	Drag	Lift	N/L
2	5.57150	0.01031	2/2
3	5.58032	0.01047	3/2
4	5.58285	0.01051	2/2
5	5.58351	0.01052	2/2

$We = 0.002$        $a = 1, G = 0, H = \alpha \sigma_p^2$

Level	Drag	Lift	N/L
2	5.56474	0.01053	2/2
3	5.57511	0.01064	2/2
4	5.57936	0.01062	2/2
5	5.58131	0.01059	2/2

$We = 1.0$        $a = 1, G = 0, H = 0$

Level	Drag	Lift	N/L
2	20.8412	0.32761	7/7
3	17.7123	0.22910	6/8
4	15.0096	0.14311	6/9
5	12.58895	0.07002	6/9

$We = 1.0$        $a = 1, G = 0, H = \alpha \sigma_p^2$

Level	Drag	Lift	N/L
2	5.03961	-0.00172	4/2
3	4.93834	-0.00210	4/3
4	4.84483	-0.00252	3/3
5	4.77541	-0.00276	3/4

**Level independent solver !**



## Robust numerical and algorithmic tools are available using

- ✓ **Monolithic Finite Element Method (M-FEM)**
- ✓ **Classical and Log Conformation Reformulation (LCR)**
- ✓ **Edge Oriented stabilization (EO-FEM)**
- ✓ **Fast Newton-Multigrid Solver with local MPSC smoother**

for the simulation of nonlinear flow with viscoelastic behaviour

## Advantages

- ✓ **No CFL-condition restriction due to the fully implicit coupling**
- ✓ **Positivity preserving**
- ✓ **Higher order and local adaptivity**

