

## FEM techniques for interfacial flows

How to avoid the explicit reconstruction of interfaces

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# Let's start with Benchmarking...



S. Hysing, S. Turek, D. Kuzmin, N. Parolini, E. Burman, S. Ganesan, and L. Tobiska; *Proposal for quantitative benchmark computations of bubble dynamics*, Submitted to Int. J. Numer. Meth. Fluids.



## Benchmark: 2D Rising Bubble

<u>Aim:</u> Development of **quantitative** two-phase flow benchmarks for rigorous evaluation of new and existing methods





Figure: Initial configuration and boundary conditions for the test cases

## Benchmark Test Cases

Test Case	1	2	
$\rho_1$ (liquid)	1000	1000	
$\rho_2$ (gas)	1	100	
$\mu_1$ (liquid)	10	10	
$\mu_2$ (gas)	0.1	1	
<i>g</i> <sub>y</sub>	-0.98	-0.98	
σ	1.96	24.5	
Re	35	35	
Eo	125	10	
$\rho_1/\rho_2$	1000	10	
$\mu_1/\mu_2$	100	10	





	Group and Affiliation	Code/Method
1	Uni. Dortmund, Inst. of Applied Math.	TP2D
	S. Turek, D. Kuzmin, S. Hysing	FEM-Level Set
2	EPFL Lausanne, Inst. of Analysis and Sci. Comp.	FreeLIFE
	E. Burman, N. Parolini	FEM-Level Set
3	Uni. Magdeburg, Inst. of Analysis and Num. Math.	MooNMD
	L. Tobiska, S. Ganesan	FEM-ALE



## Visual comparison: Test Case 1





## Benchmark quantities - circularity



## Benchmark quantities - center of mass



## Benchmark quantities - rise velocity



## Visual comparison: Test Case 2





## Benchmark quantities - circularity



## Benchmark quantities - center of mass



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## Benchmark quantities - rise velocity





### Conclusions

- Proposed two benchmarks for rigorous evaluation
- Established target reference values for the first benchmark
- Hinted at difficulties during break up in the second benchmark
- Look at www.featflow.de for participation

Interfacial flows: Challenging even in 2D!





# Interfacial flows



# Numerical simulation of interfacial or two-phase flows with immiscible fluids poses some challenging problems

### Issues

- Accurate interface tracking
- Mass conservation
- Resolution of discontinuous fluid properties
- Treatment of interfacial boundary conditions

Numerical simulation of interfacial or two-phase flows with immiscible fluids poses some challenging problems

Issues

# How to devise an efficient solution algorithm which addresses these issues?

• The Navier-Stokes equations govern incompressible fluid flow

$$\rho(\mathbf{x}) \left( \frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla) \mathbf{u} \right) = -\nabla \rho + \nabla \cdot \left( \mu(\mathbf{x}) (\nabla \mathbf{u} + (\nabla \mathbf{u})^T) \right) + \rho(\mathbf{x}) \mathbf{g}$$
$$\nabla \cdot \mathbf{u} = 0$$

with varying density  $\rho(\mathbf{x})$  and viscosity  $\mu(\mathbf{x})$  fields

nterfacial Boundary Conditions

• Direct interface conditions

$$[\mathbf{u}]|_{\Gamma} = \mathbf{0}, \quad -\left[-\rho\mathbf{I} + \mu(\nabla\mathbf{u} + (\nabla\mathbf{u})^{T})\right]|_{\Gamma} \cdot \hat{\mathbf{n}} = \sigma\kappa\hat{\mathbf{n}}$$

Implicit conditions by weighted volume forces

$$\mathbf{f}_{st}|_{\Gamma} = \sigma \kappa \hat{\mathbf{n}}$$

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Interfacial Boundary Conditions

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• Implicit conditions by weighted volume forces

$$\mathbf{f}_{st}|_{\Gamma} = \sigma \kappa \hat{\mathbf{n}}$$

## Saddle-point systems after Discretization

Given  $\mathbf{u}^n$  and time step  $k = \Delta t$ , then solve for  $\mathbf{u} = \mathbf{u}^{n+1}$  and  $p = p^{n+1}$ 

$$\begin{cases} S\mathbf{u} + kBp = \mathbf{f}^{n+1} \\ B^{\mathsf{T}}\mathbf{u} = 0 \end{cases}$$

with

$$S\mathbf{u} = [M_{\rho} + \theta k N(\mathbf{u})]\mathbf{u}$$

$$N(\mathbf{v}) = -
abla \cdot \left( \mu(\mathbf{x}) (
abla \mathbf{u} + (
abla \mathbf{u})^T) \right) + 
ho(\mathbf{x}) (\mathbf{u} \cdot 
abla) \mathbf{v}$$

B and  $B^T$  are discrete counterparts to the grad and div operators





### We prefer to use an operator splitting approach:

Given  $\mathbf{u}^n$  and  $p^n$ , then do for time level n + 1:

- Solve the Navier-Stokes equations, with the interface given at Γ<sup>n</sup>, to obtain u<sup>n+1</sup> and p<sup>n+1</sup>
- Optionally perform grid adaptation)
- Move the interface according to the given physics
- Perform post-processing; e.g. computation of normals and curvature
- Optionally perform grid adaptation)
- (Re-iterate time loop if deemed necessary)





# Interface Tracking





The tasks of an interface tracking algorithm are:

- Correctly propagate the interface given a velocity field
- Minimize additional mass loss due to the algorithm itself
- Enable easy and quick identification of both the interface and the different phases
- Enable easy reconstruction of interface normal vectors and curvature
- Be relatively easy to implement and efficient to solve
- Ideally treat coalescence and break-up automatically

## Interface Tracking

Two main schools of thought...

Lagrangian



### Eulerian





## Interface Tracking

Which method should we use then?

- Pure Lagrangian Approach
- Front Tracking Method
- Segment Projection Method
- Marker and Cell (MAC)
- Volume of Fluid (VOF)
- Phase Field Method
- Level Set Method (LS)
- Particle Level Set Method (PLS)
- Combination (LS-VOF)





Our selection criteria

- Eulerian, simplifies implementation
- Allow for discretization and solution with fast and efficient PDE solver techniques
- Handle coalescence and break-up automatically
- Allow for **global** reconstruction of normals and curvature



## Interface Tracking

Our selection criteria

## Level Set Method



## Interface Tracking

### The Level Set Method

The general idea is to embed an interface in a higher dimensional function  $\phi$ . The interface movement is then governed by a standard transport equation

$$\frac{\partial \phi}{\partial t} + \mathbf{v} \cdot \nabla \phi = \mathbf{0}$$

with initial condition  $\phi(\mathbf{x} = \Gamma, t = 0) = 0$ 







The Level Set Method - Properties

- A smoothness constraint on the level set field relaxes the requirements on the solver
- The natural choice is to restrict the level set field to be a signed distance function

$$|
abla \phi| = 1$$

• At any given point the magnitude of the level set function will thus represent the shortest distance to the interface





### The Level Set Method - Advantages

- The smooth LS function allows for higher order discretization techniques to be employed
- The governing transport equation can be solved with efficient standard solvers
- The interface is implicitly but also exactly defined
- Break up and coalescence is treated automatically
- Geometrical quantities such as normals and curvature can be reconstructed globally





## Problems

### The Level Set Method - Problems

• The velocity field **v** convecting the level set function can not in general be expected to maintain  $\phi$  as a distance function

$$\frac{\partial \phi}{\partial t} + \mathbf{v} \cdot \nabla \phi = \mathbf{0}$$

- ${\ensuremath{\, \bullet }}$  The distance function property is only preserved if  $\nabla {\mathbf v} \cdot \nabla \phi = {\mathbf 0}$
- Stretching and folding can cause eventual solver failure





## Problems

The Level Set Method - Problems

## Periodic Reinitialization



Several methods exist for constructing a distance function from given interface data. The simplest ones use interface approximations for the reconstruction

- Redistancing via "Brute Force"
- Algebraic Newton approach

The more complex algorithms deal with solving the Eikonal equation  $|\nabla \phi| = F$  on a fixed mesh

- The Fast Marching Method
- The Fast Sweeping Method
- PDE Based redistancing
- Branch and Bound approach to redistancing



## **Reinitialization Methods**

The different methods naturally have their respective strengths and weaknesses

- Brute force method
  - + Very robust
  - Scales quite badly,  $\mathcal{O}(N * M)$
- Algebraic Newton Approach
  - + Potentially very fast,  $\mathcal{O}(N)$ 
    - Convergence dependent upon approximate distance function
- The Fast Marching Method
  - + Easy to implement in an unstructured context
  - Needs a heap structure to sort marching order,  $+\mathcal{O}(logN)$
- The Fast Sweeping Method
  - + Potentially very fast,  $\mathcal{O}(N)$ 
    - Difficult to generalize to fully unstructured grids





• The stationary limit of the following PDE can also be used to apply reinitialization to a given approximate distance field  $\phi_0$ 

$$rac{\partial \phi^*}{\partial t} + \mathbf{q} \cdot 
abla \phi^* = S(\phi_0), \qquad \mathbf{q} = S(\phi_0) rac{
abla \phi^*}{|
abla \phi^*|}$$

- where  $S(\phi_0)$  is an appropriately chosen sign function
- no local reconstruction necessary, but very expensive !





Test Case

A non-trivial test case was chosen to represent a typical level set computation with the following properties

- Include both smooth regions and shocks
- Exact solution available:

$$\phi(\mathbf{x}) = \min(2.25 - y, \sqrt{x^2 + (y - 1)^2} - 0.4)$$



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## Results - CPU Time



## **Results** - Accuracy







### Reinitialization

- The fast marching method scales very well with increasing grid density
- Fast marching is very fast even for very dense grids ( $O(10^6)$  grid points)
- The second order update only costs marginally more than the first order version
- The fast marching method is therefore our preferred algorithm! (but: explicit local reconstruction)





# Treatment of Surface Tension





- Interfacial or two-phase flow where capillary surface tension forces are dominant poses some challenging problems
- Surface tension effects are generally modeled both explicitly in time and space leading to the capillary time step restriction

$$\Delta t_{num}^{(ca)} < \sqrt{rac{\langle 
ho 
angle \, h^3}{2\pi\sigma}}$$

### Goal

Remove the capillary time step constraint while retaining a fully Eulerian interface description



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## Computation of Surface Tension

Surface tension effects can essentially be included in the Navier-Stokes equations in two different ways

• Explicit interface reconstruction and direct evaluation

$$\mathbf{f}_{st} = \int_{\Gamma} \sigma \kappa \hat{\mathbf{n}} \ d\Gamma$$

• Implicit incorporation via the continuum surface force (CSF) model

$$\mathbf{f}_{st} = \int_{\Omega} \sigma \kappa \hat{\mathbf{n}} \delta(\Gamma) \ d\Omega$$





## Definitions

Definition (Tangential gradient)

The tangential gradient of a function f, which is differentiable in an open neighborhood of  $\Gamma,$  is defined by

 $\underline{\nabla} f(x) = \nabla f(x) - (\hat{\mathbf{n}}(x) \cdot \nabla f(x))\hat{\mathbf{n}}(x), \quad x \in \Gamma$ 

where  $\nabla$  denotes the usual gradient in  $\mathbb{R}^d$ 

Definition (Laplace-Beltrami operator)

If f is two times differentiable in a neighborhood of  $\Gamma$ , then we define the Laplace-Beltrami operator of f as

 $\underline{\Delta}f(x) = \underline{\nabla} \cdot (\underline{\nabla}f(x)), \quad x \in \Gamma$ 

## **Definitions and Derivation**

### Theorem

A theorem of differential geometry states that

$$\underline{\Delta}$$
id <sub>$\Gamma$</sub>  =  $\kappa \hat{\mathbf{n}}$ 

where  $\kappa$  is the mean curvature and  $\mathrm{id}_\Gamma$  is the identity mapping on  $\Gamma$ 

### Derivation

First take surface tension force source term, multiply it with the test function space  $\bm{v},$  and apply partial integration

$$\begin{aligned} \mathbf{f}_{st} &= \int_{\Gamma} \sigma \kappa \hat{\mathbf{n}} \cdot \mathbf{v} \ d\Gamma &= \int_{\Gamma} \sigma (\underline{\Delta} \mathrm{id}_{\Gamma}) \cdot \mathbf{v} \ d\Gamma &= \\ &= -\int_{\Gamma} \sigma \underline{\nabla} \mathrm{id}_{\Gamma} \cdot \underline{\nabla} \mathbf{v} \ d\Gamma + \int_{\gamma} \sigma \partial_{\gamma} \mathrm{id}_{\Gamma} \cdot \mathbf{v} \ d\gamma \end{aligned}$$





## Fully Implicit Evaluation in Space

• Boundary integrals can be transformed to volume integrals with the help of a Dirac delta function  $\delta(\Gamma, \mathbf{x})$ 

$$\begin{aligned} \mathbf{f}_{st} &= \int_{\Omega} \sigma \kappa \hat{\mathbf{n}} \cdot \mathbf{v} \ \delta(\Gamma, \mathbf{x}) \ d\mathbf{x} &= \int_{\Omega} \sigma(\underline{\Delta} \mathrm{id}_{\Gamma}) \cdot (\mathbf{v} \delta(\Gamma, \mathbf{x})) \ d\mathbf{x} \\ &= -\int_{\Omega} \sigma \underline{\nabla} \mathrm{id}_{\Gamma} \cdot \underline{\nabla} (\mathbf{v} \delta(\Gamma, \mathbf{x})) \ d\mathbf{x} \ = -\int_{\Omega} \sigma \underline{\nabla} \mathrm{id}_{\Gamma} \cdot \underline{\nabla} \mathbf{v} \ \delta(\Gamma, \mathbf{x}) \ d\mathbf{x} \end{aligned}$$

• Application of the semi-implicit time integration

$$\Gamma^{n+1} = \Gamma^n + \Delta t \, \mathbf{u}^{n+1}$$

yields

$$\begin{aligned} \mathbf{f}_{st} &= -\int_{\Omega} \sigma \underline{\nabla} (\mathrm{id}_{\Gamma})^{n} \cdot \underline{\nabla} \mathbf{v} \ \delta(\Gamma^{n}, \mathbf{x}) \ d\mathbf{x} \\ &- \Delta t^{n+1} \int_{\Omega} \sigma \underline{\nabla} \mathbf{u}^{n+1} \cdot \underline{\nabla} \mathbf{v} \ \delta(\Gamma^{n}, \mathbf{x}) \ d\mathbf{x} \end{aligned}$$





## Regularization

• Regularization of  $\delta(\Gamma, \mathbf{x})$  can easily be accomplished with the help of a distance function

$$\delta_{\boldsymbol{\epsilon}}(\boldsymbol{\Gamma}, \mathbf{x}) = \delta_{\boldsymbol{\epsilon}}(dist(\boldsymbol{\Gamma}, \mathbf{x})),$$

where  $dist(\Gamma, \mathbf{x})$  gives the minimum distance from  $\mathbf{x}$  to  $\Gamma$ 

 ${\, \bullet \,}$  The regularized continuous delta function  $\delta_\epsilon$  is defined as

$$\delta_{\epsilon}(x) = \begin{cases} \frac{1}{\epsilon}\varphi(x/\epsilon) & |x| \leq \epsilon & = \ mh, \\ 0 & |x| > \epsilon & = \ mh, \end{cases}$$

where h is the mesh spacing which together with the constant m defines the support  $\epsilon$  of the regularized delta function,  $\varphi$  is a characteristic function determining the kernel shape

 A second order method can alternatively be constructed according to [Engquist, Tornberg, and Tsai: J. Comput. Phys. 207:28-51, 2005]

## Regularization

There are several possible choices for  $\varphi(\xi)$ , common choices are:

- The linear hat function  $arphi^1(\xi) = 1 |\xi|$
- The commonly used cosine approximation  $\varphi^2(\xi) = \frac{1}{2}(1 + \cos(\frac{\pi\xi}{2}))$
- Higher order polynomials, for example;

$$\varphi^{3}(\xi) = \frac{312}{512} (3 - 20\xi^{2} + 42\xi^{4} - 36\xi^{6} + 11\xi^{8})$$





## Implicit Surface Tension Force Expression

### The surface tension forces are finally given by ...

Implicit Surface Tension Force Expression

$$\begin{split} \mathbf{f}_{st} &= -\int_{\Omega} \sigma \ \delta_{\boldsymbol{\epsilon}}(dist(\boldsymbol{\Gamma}^{n},\mathbf{x})) \ \underline{\nabla}(\widetilde{\mathrm{id}}_{\boldsymbol{\Gamma}})^{n} \cdot \underline{\nabla} \mathbf{v} \ d\mathbf{x} \\ &- \Delta t^{n+1} \int_{\Omega} \sigma \ \delta_{\boldsymbol{\epsilon}}(dist(\boldsymbol{\Gamma}^{n},\mathbf{x})) \ \underline{\nabla} \mathbf{u}^{n+1} \cdot \underline{\nabla} \mathbf{v} \ d\mathbf{x} \end{split}$$





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A number of key points makes the *level set method* an ideal candidate for interface tracking algorithm when implementing the proposed surface tension force expressions

- Distance functions are in general readily available allowing for simple construction of the regularized Dirac delta functions
- Geometrical quantities such as normal and tangent vectors can be reconstructed globally, eliminating the need to extend these quantities from the interface separately
- The level set method can be coupled with the finite element method giving access to the variational form of the equations





### Method Validation

### Numerical Examples



## Validation, Laplace-Young Law

A perfect circular static bubble should follow the Laplace-Young law

$$p_{inside} = p_{outside} + \frac{\sigma}{r}$$



Figure: Pressure cut-line for four different mesh sizes.



## Example, Oscillating Bubble (CSF)



Figure: Evolution of an oscillating bubble; standard explicit CSF method.

## Example, Oscillating Bubble (CSF-LBI)



Figure: Evolution of an oscillating bubble; semi-implicit CSF-LBI method.

## Example, Rising Bubble (CSF)



Figure: Evolution of a rising bubble with the standard explicit CSF method.

## Example, Rising Bubble (CSF-LBI)



Figure: Evolution of a rising bubble with the semi-implicit CSF-LBI method.



A new implicit surface tension variant has been proposed which relaxes the capillary time step restriction imposed on explicit implementations

Additional advantages

- Fully implicit in space
- Is easily implemented when using the level set method together with finite elements
- Explicit computation of curvature not necessary
- Conceptually identical algorithm in 3D





# Grid Deformation Techniques





## Grid Deformation

### Monitor function

• In the case of interfacial flow, then the **monitor function** f can be constructed from a distance function, giving the shortest distance to the interface, and possibly also weighting the interface curvature  $\kappa$ 

 $f = f(|\phi(x)|, \kappa(x))$ 

- This makes the *level set method* ideally suited to use as interface tracking algorithm, since both the **distance function and curvature** are defined **globally**
- Special high-performance-computing techniques and hardware (GPU computing, Cell processor, Sony PSP3) can be exploited

## **Grid** Deformation







## Static Bubble Test



Figure: Pressure field for a static bubble with and without grid deformation

## Static Bubble Test





Figure: Velocity error for a static bubble with and without grid deformation

GridLevel	3	4	5	6			
Tensor product grid							
U error	$3.7 \cdot 10^{-2}$	$1.1 \cdot 10^{-2}$	$1.0 \cdot 10^{-2}$	$5.6 \cdot 10^{-3}$			
P error	5.583%	1.433%	0.212%	0.037%			
Adapted grid							
U error	$2.0 \cdot 10^{-2}$	$7.3 \cdot 10^{-3}$	$3.5 \cdot 10^{-3}$	$1.8 \cdot 10^{-3}$			
P error	2.969%	0.261%	0.042%	0.013%			





## Future Research Directions

### In development...

- ALE techniques coupled with time dependent grid deformation
- Inclusion of *pressure separation* techniques to improve the velocity and pressure
- Linear high order *edge stabilization* for convection of the level set field
- $\mathbb{Q}_2\mathbb{P}_1$  finite element approximation for the NS-equations and  $\mathbb{Q}_n$  for the level set equation
- Contact angle, heat transfer, and solidification effects
- 3D + benchmarking